



# Observational Effects of Banded Repeating FRBs

Kshitij Aggarwal<sup>1,2</sup> <sup>1</sup> West Virginia University, Department of Physics and Astronomy, P.O. Box 6315, Morgantown, WV, USA; [ka0064@mix.wvu.edu](mailto:ka0064@mix.wvu.edu)<sup>2</sup> Center for Gravitational Waves and Cosmology, West Virginia University, Chestnut Ridge Research Building, Morgantown, WV, USA

Received 2021 August 10; revised 2021 September 24; accepted 2021 September 26; published 2021 October 8

## Abstract

Recent observations have shown that repeating fast radio bursts (FRBs) exhibit band-limited emission, whose frequency-dependent amplitude can be modeled using a Gaussian function. In this analysis, we show that banded emission of FRBs can lead to incompleteness across the observing band. This biases the detected sample of bursts and can explain the various shapes of cumulative energy distributions seen for repeating FRBs. We assume a Gaussian shape of the burst spectra and use simulations to demonstrate the above bias using an FRB 121102-like example. We recovered energy distributions that showed a break in power law and flattening of power law at low energies, based on the fluence threshold of the observations. We provide recommendations for single-pulse searches and analysis of repeating FRBs to account for this incompleteness. Primarily, we recommend that burst spectra should be modeled to estimate the intrinsic fluence and bandwidth of the burst robustly. Also, bursts that lie mainly within the observing band should be used for analyses of energy distributions. We show that the bimodality reported in the distribution of energies of FRB 121102 by Li et al. disappears when burst bandwidth, instead of the center frequency of the observation, is used to estimate energy. Subbanded searches will also aid in detecting band-limited bursts. All the analysis scripts used in this work are available in a Github repository ([https://github.com/KshitijAggarwal/banded\\_repeater\\_analysis](https://github.com/KshitijAggarwal/banded_repeater_analysis)).

*Unified Astronomy Thesaurus concepts:* [Radio transient sources \(2008\)](#); [Extragalactic radio sources \(508\)](#); [Radio bursts \(1339\)](#); [Astronomy data analysis \(1858\)](#)

## 1. Introduction

Fast radio bursts (FRBs) show a wide variety of properties (Petroff et al. 2019). There is an ongoing effort to disentangle the properties that arise from the FRB source itself (i.e., intrinsic) as well as those properties introduced due to various selection effects. FRBs that emit multiple bursts, called repeaters, appear to be distinct in their observational characteristics as compared to the apparent nonrepeaters (The CHIME/FRB Collaboration et al. 2021). Two properties associated primarily with repeating FRBs are subburst drifting and band-limited emission (Law et al. 2017; Hessels et al. 2019). In this Letter, the effect we focus on is the band-limited emission of repeaters and we discuss the various observational biases caused by it.

The source of the band-limited nature of FRB emission is currently not understood. However, there exist multiple progenitor and propagation models that try to explain the band-limited nature of FRBs (Cordes et al. 2017; Lu & Kumar 2018; Metzger et al. 2019; Beniamini & Kumar 2020; Simard & Ravi 2020). Further, it has been reported that the emission of the first repeating FRB, FRB 121102, favors 1600 MHz, and there is a lack of emission observed below 1200 MHz (Aggarwal et al. 2021a; Platts et al. 2021). It is unclear whether the emission behavior changes below that frequency and if it is below the detection threshold or if it is just not present (Platts et al. 2021). Similar banded<sup>3</sup> emission has also been reported for other repeating FRBs (Pastor-Marazuela et al. 2021; Aggarwal et al. 2021a; Li et al. 2021). Further, it has been reported that the peak emission frequency appears to be random. Also, ultra-wideband observations have shown that there is no evidence for oscillations in the spectra, i.e., the burst emission is present only within the narrow envelope and is not simultaneously present at any other frequency (Law et al. 2017;

Kumar et al. 2021). The periodicity in burst activity and its frequency dependence further complicates the interpretation of this effect (Aggarwal et al. 2020; Pastor-Marazuela et al. 2021).

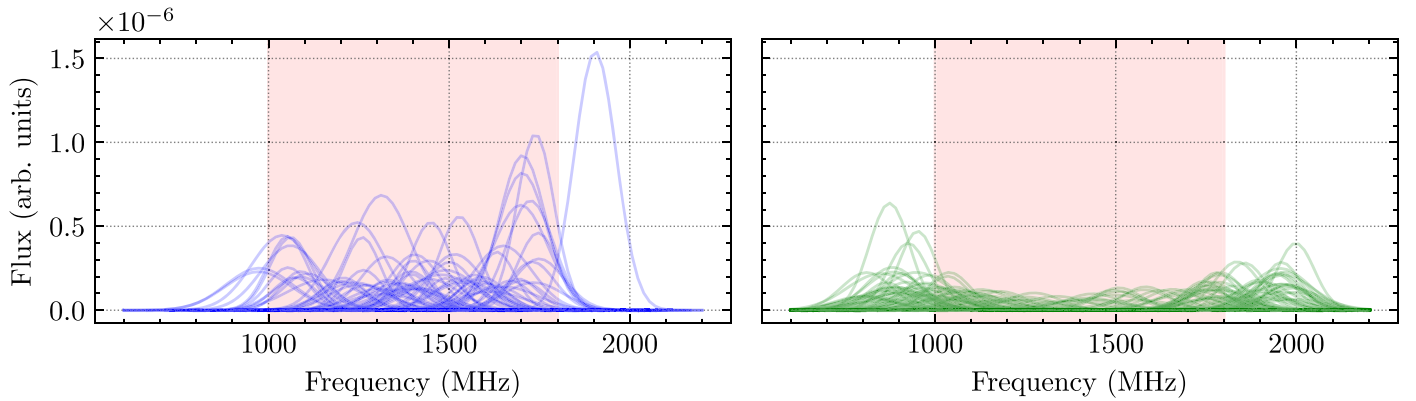
Energy distribution of repeating FRBs also shows a variety of shapes and features: a simple power law, a broken power law, a smooth flattening of the power law at low energies, etc. (Cruces et al. 2020; Pastor-Marazuela et al. 2021; Aggarwal et al. 2021a; Li et al. 2021). The shape of the energy distribution can provide important information about the emission mechanism of the FRB source. Giant pulses from neutron stars and high energy emission from magnetars has been described using a power-law distribution (Bera & Chengalur 2019; Cheng et al. 2020). A power-law index for FRBs similar to that seen from neutron stars might imply a common origin. In some studies of repeaters, the deviation of energy distribution from a simple power law has been attributed to the intrinsic emission process of the FRB. But, in this Letter, we demonstrate that many of the observed shapes of the energy distributions, can be attributed to biases due to the band-limited nature of emission of repeaters. It is therefore necessary to account for these biases, before making conclusions about the intrinsic nature of FRB emission.

We start by discussing our modeling methods in Section 2, followed by results demonstrated using an example in Section 3, and we discuss the implications in Section 4. We then conclude with some recommendations for single-pulse search and analysis of repeater bursts in Section 5.

## 2. Methods

This section briefly discusses the methods we used to model the FRB spectra and various observational effects. We follow a

<sup>3</sup> We use “band-limited” and “banded” emission interchangeably throughout this Letter. Both of these refer to the finite bandwidth emission of the bursts.



**Figure 1.** Simulated spectra modeled using a Gaussian function. The mean and standard deviation have been sampled from a uniform distribution. The energy distribution is assumed to be a power law with a slope of  $-1.8$ . The left panel shows the detected spectra, while the right panel shows the ones that were not detected, using a constant fluence threshold. The observing band is shown in red. Spectra that had enough energy within the band were detected, while the ones without enough signal within the observing band were not detected. See Section 2.2 for more details.

two-step approach: (1) simulate a population of FRBs with spectral properties described by predefined functions, and (2) estimate the detectable FRBs from this sample based on a sensitivity threshold.

### 2.1. Generating a Population of Repeater Bursts

Previous studies have shown that the spectra of repeating FRBs can be modeled using a Gaussian function (Law et al. 2017; Aggarwal et al. 2021a; Kumar et al. 2021; Pleunis et al. 2021). Therefore, we assume that the repeater burst’s spectra follow a Gaussian function, parameterized by a mean ( $\mu_f$ ) and a standard deviation ( $\sigma_f$ ). We also assume that  $\mu_f$  and  $\sigma_f$  follow a predefined distribution (Gaussian or uniform). Finally, we assume that the cumulative energy distribution follows a power law with a slope ( $\alpha$ ). We then draw 50,000 samples from the above three distributions each to represent 50,000 bursts from a repeater. We convert the energy into an intrinsic fluence using its spectral bandwidth and assuming a nominal distance to the source.

### 2.2. Applying Selection Effects

Next, we aim to determine the bursts that a given observational system will detect. To simplify the analysis, we ignore the effect of any signal lost due to the nonidealized dispersion measure (DM) and boxcar search step (see Keane & Petroff 2015; Agarwal et al. 2020; Aggarwal et al. 2021a, 2021b, on details of those effects), and only consider two selection effects: observational bandwidth and fluence limit. Due to the limited observing bandwidth of most instruments and the distribution of burst spectra across a wide frequency range (see Figure 1), the observed fluence of the burst will depend on the fraction of the burst spectra that lies within the observational band. Therefore, for each simulated burst, we estimate its observed fluence by integrating the spectra (using the bursts’  $\mu_f$  and  $\sigma_f$ ) within the limits of observing bandwidth. Therefore, assuming Gaussian spectra of the form,

$$\mathcal{G}(\nu; \mu_\nu, \sigma_\nu) = \frac{1}{\sigma_\nu \sqrt{2\pi}} \exp\left(-\frac{1}{2} \frac{(\nu - \mu_\nu)^2}{\sigma_\nu^2}\right) \quad (1)$$

where  $\nu$  is the observing frequency. The observed fluence is given by:

$$S_{\text{obs}} = S_{\text{int}} \int_{\nu_{\text{start}}}^{\nu_{\text{end}}} \mathcal{G}(\nu; \mu_\nu, \sigma_\nu) d\nu. \quad (2)$$

Where  $S_{\text{int}}$  and  $S_{\text{obs}}$  are the intrinsic and observed burst fluence,  $\nu_{\text{start}}$  and  $\nu_{\text{end}}$  are the start and end frequencies of the observing band. If the burst spectra lie primarily within the bandwidth, then the integral in the above equation will be close to one, and so the observed fluence will be very similar to the intrinsic fluence. As the fraction of the burst signal outside the band increases, the observed fluence will lessen as compared to intrinsic fluence.

If this observed fluence is greater than the fluence threshold of the search pipeline, then the burst will be detected. Therefore, out of the population of bursts that do not lie primarily within our observing band, we are sensitive to detecting only the bright ones (as illustrated in Figure 1). This will introduce strong selection effects in the properties of the detected bursts.

### 2.3. Estimating Fluence and Bandwidth

Typically, the detected bursts’ fluence (and then energy) is determined only from the signal visible in the observing band. This is done using the signal-to-noise ratio (S/N) obtained from the dedispersed and frequency averaged time-series profile of the burst. However, this procedure underestimates the intrinsic fluence of the burst, as it is sensitive to the signal present only in the observing band. On the other hand, if we model the observed spectra using a Gaussian function, then the total fluence (including the signal not visible in the band) can be estimated and determined (see Section 3.5 of Aggarwal et al. 2021a). Using this burst modeling, it is possible to obtain a more robust estimate of the intrinsic fluence of the burst (i.e.,  $S_{\text{int}}$ ).

Similarly, the bandwidth of the bursts is typically determined by manually identifying the range of frequency channels in which the burst signal is visible. For band-limited bursts that lie on the edge of the band, this would lead to underestimation of the burst bandwidths. Again, suppose we model the burst spectra using a Gaussian model. In that case, we can accurately determine the full width at half-maximum (FWHM) of all the

bursts robustly, even if only a small fraction of burst spectra is within the observing band.

We use the following equation to estimate the isotropic energy of the bursts from the fluence ( $S$ ), bandwidth (FWHM), and distance ( $D_L$ ) of the bursts

$$E = 4\pi 10^{-23} \left( \frac{D_L}{\text{cm}} \right)^2 \left( \frac{S}{\text{Jy s}} \right) \left( \frac{\text{FWHM}}{\text{Hz}} \right) \text{erg}. \quad (3)$$

Henceforth, we refer to the S/N derived fluence as  $F_{S/N}$ , and energies estimated using  $F_{S/N}$  and manually identified bandwidth as  $E_{S/N}$ . We will use  $F_{\text{fit}}$  and  $E_{\text{fit}}$  when the fluence and energy are determined using fitting.

The above three steps give us a sample of simulated and detected repeater bursts that can now be compared to infer the observational biases and incompleteness, with some assumptions on the intrinsic ( $\mu_f$ ,  $\sigma_f$ , and  $\alpha$ ) and observational (fluence threshold and bandwidth) parameters. This is explored in the next section.

### 3. Results

Here we consider a simple example to report some of the observed effects of the band-limited nature of burst spectra. We consider an FRB 121102-like repeater observed at varying detection thresholds. We assume the following intrinsic properties for this repeater: cumulative energy distribution follows a single power law with a fixed slope ( $\alpha = -1.5$ ), a distance of 972 Mpc, normal distribution of  $\mu_f$  with mean 1650 MHz and standard deviation of 250 MHz, and normal distribution of  $\sigma_f$  with mean of 300 MHz and standard deviation of 250 MHz. We also assume that our observing bandwidth is 800 MHz with a center frequency of 1375 MHz. The choice of these values is inspired by the observed properties of FRB 121102 reported by Aggarwal et al. (2021a). We then observe this repeater at three different fluence limits (i.e., sensitivity limit of the observing system) of 0.02, 0.1, and 0.4 Jy ms. As discussed previously, some of these bursts will not be detected by the search system, as not enough burst signal is present within the band. In the following subsections, we discuss three unique observational effects observed for this simulated repeater. We also try to explain the various observed properties of the two most studied repeaters so far, FRB 121102 and FRB 180916, in the context of the incompleteness due to the banded nature of their bursts.

#### 3.1. Cumulative Energy Distribution

Many different shapes and slopes have been reported for the cumulative energy distribution of repeater burst energies (Law et al. 2017; Cruces et al. 2020; Pastor-Marazuela et al. 2021; Aggarwal et al. 2021a). Here we discuss how telescope sensitivity plays a crucial role in determining the observed shape of the cumulative energy distribution by influencing the bursts that are detected by the system.

The bottom row of Figure 2 shows the cumulative energy distribution of bursts detected with two different fluence thresholds. Red pluses show the intrinsic power law, green crosses show the  $E_{S/N}$  for the detected bursts, and blue triangles represent the  $E_{\text{fit}}$  of the detected bursts. The trend of blue triangles and green crosses deviates significantly from a simple intrinsic power law due to the absence of weak bursts that are not detected. Notably, the shape of observed cumulative energy distribution changes with sensitivity thresholds. As the sensitivity of the observations decreases (left to right in Figure 2), the observed distributions

deviate from a single power law to a broken power law, which appears as a smooth turn at even lower thresholds. The broken power law is similar to what has been seen for FRB 121102 (Cruces et al. 2020; Aggarwal et al. 2021a) while the smooth turn in power law was reported for FRB 180916 (Pastor-Marazuela et al. 2021). Here, we have shown that both of these effects can occur due to the band-limited nature of repeater bursts.

#### 3.1.1. Challenges

Various possible shapes of the observed cumulative distribution make it challenging to interpret and infer intrinsic FRB properties. A break in the power law, if present, might indicate the real completeness limit of the system. This can be observed by the green crosses in the bottom left plot of Figure 2. The higher energy slope appears similar to the intrinsic one. Therefore, a break in observed power law might indicate that the higher energy power law follows the intrinsic shape and may be used to draw inferences about the intrinsic properties of the FRB. A smooth turnover in cumulative distribution (right panel of Figure 2), on the other hand, renders a completeness limit derived from such an analysis inaccurate, and even the higher energy slope might not represent the intrinsic distribution of energies. Even the distribution of  $E_{\text{fit}}$  might not reflect the real power law, as the detection threshold would also bias it.

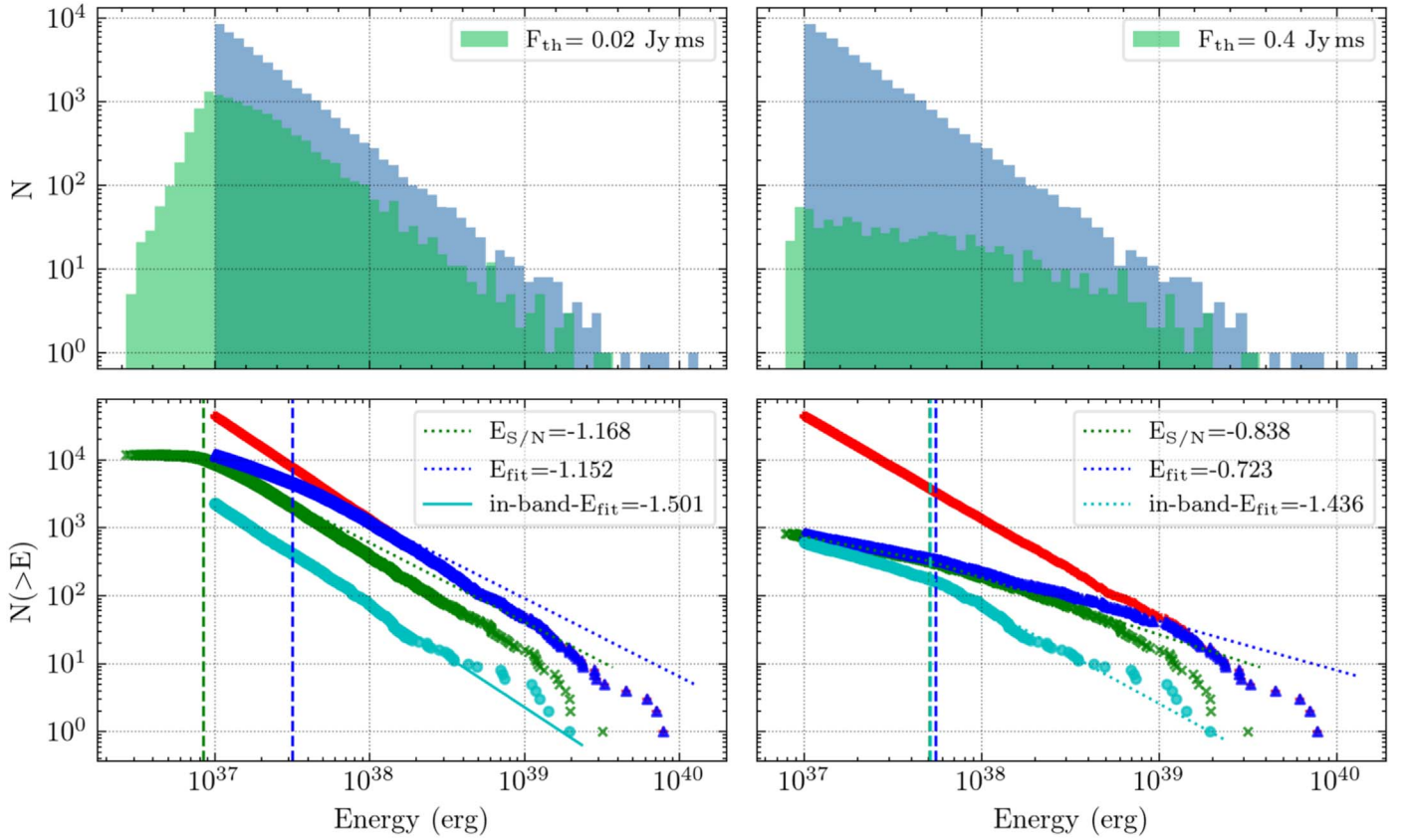
#### 3.1.2. Using In-band Bursts

When using  $E_{\text{fit}}$ , although we are using the correct fluences of the bursts, we still would have missed a population of weak bursts that were not primarily within our observing band (see Section 2.2). Therefore, we would detect a larger sample of bright bursts, which will make the energy distribution flatter. We can account for this by analyzing only the in-band bursts, i.e., the bursts whose spectra lie primarily within our observing band.<sup>4</sup> These bursts are labeled as in-band- $E_{\text{fit}}$  and shown as cyan circles in Figure 2. These in-band bursts provide a more robust and reliable estimate of the distribution of burst energies, even when the observations are not very sensitive. This is because our observations are complete to the in-band bursts, and by using fitting to determine burst properties, we have mitigated both selection effects listed earlier (see Section 2.2). The slope of the in-band burst energies follows the intrinsic distribution. But, if the distribution shows a turnover or flattening, then even this method cannot be used to reliably estimate the intrinsic power-law slope of the bursts.

#### 3.2. Energy Distribution

Continuing the previous example, the top row in Figure 2 shows the histogram of  $E_{S/N}$  for two fluence thresholds. Again, the intrinsic cumulative distribution of burst energies follows a power law. The detected energy distribution again changes significantly with the fluence threshold, and even in this simple case, it shows a variety of shapes. All of these shapes can be attributed to the missed bursts and inaccurate estimates of burst energies. Using  $E_{\text{fit}}$  leads to distributions that are more similar to the intrinsic one, and using  $E_{\text{fit}}$  for only in-band bursts would be an even more accurate representation of intrinsic distribution.

<sup>4</sup> A similar condition was also used by Aggarwal et al. (2021a) to select the bursts for cumulative energy distribution analysis.



**Figure 2.** Energy distributions of bursts detected at varying fluence thresholds. Different columns show different fluence thresholds—left: 0.02 Jy ms; right: 0.4 Jy ms. Top: histogram of burst energies. Intrinsic energies of the bursts are shown in blue and show the single power law, while those of detected bursts are shown in green. In this case, the burst energies are estimated using  $F_{S/N}$  (i.e., using fluence derived from signal-to-noise of the burst; see Sections 3.1 and 3.2). Bottom: cumulative energy distribution of bursts. Red pluses show the intrinsic energy distribution that follows a power law with a slope of  $-1.5$ . The other three colors show detected bursts. Green crosses represent energies estimated using  $F_{S/N}$ . Blue triangles and cyan circles are with energies estimated using fitting. Blue triangles show all the detected bursts, while cyan circles only show bursts that were primarily within the observing band. Solid lines show a single power-law fit, and dotted lines show a broken power-law fit. The vertical dashed lines show the break energy for the broken power-law fit. Values in the legends report the fitted slope for the single power-law fit, and high energy slope for broken power-law fit. As we can see, due to inaccurate energy estimation and missed bursts, the energy distributions can deviate significantly from the intrinsic distribution. See Section 3.2 for more details.

### 3.3. Distribution of Spectral Parameters

We consider two choices of intrinsic distributions for both  $\mu_f$  and  $\sigma_f$ , uniform and normal, and observed the respective distributions of the detected bursts (see Figure 3). The  $\mu_f$  distribution of detected bursts follows a normal-like distribution for both intrinsic distributions.  $\mu_f$  distribution for FRB 121102 has been reported to be an asymmetric normal distribution, with a negative skew, i.e., with a tail toward lower frequencies being drawn out (Aggarwal et al. 2021a). The FRB emission also prefers higher frequencies in the 1.4 GHz observations. This cannot be recovered using a uniform intrinsic distribution of  $\mu_f$ , as in this case, the recovered distribution would peak at the center of the observing band (bottom left panel in Figure 3). On the other hand, if the intrinsic  $\mu_f$  distribution is normal, with its mean present toward the top of the center frequency of the observing band, then an asymmetric normal distribution is recovered. The recovered  $\mu_f$  distribution of our simple example also shows a negatively skewed distribution (top left panel in Figure 3). This is because fewer bursts will be detected toward the higher part of the band. In this case, the peak of the recovered distribution will lie close to the peak of the intrinsic distribution. Hence, based on results presented in Aggarwal et al. (2021a), we can infer that the

intrinsic  $\mu_f$  distribution of FRB 121102 bursts could be normal, with a mean of  $\sim 1600$  MHz.

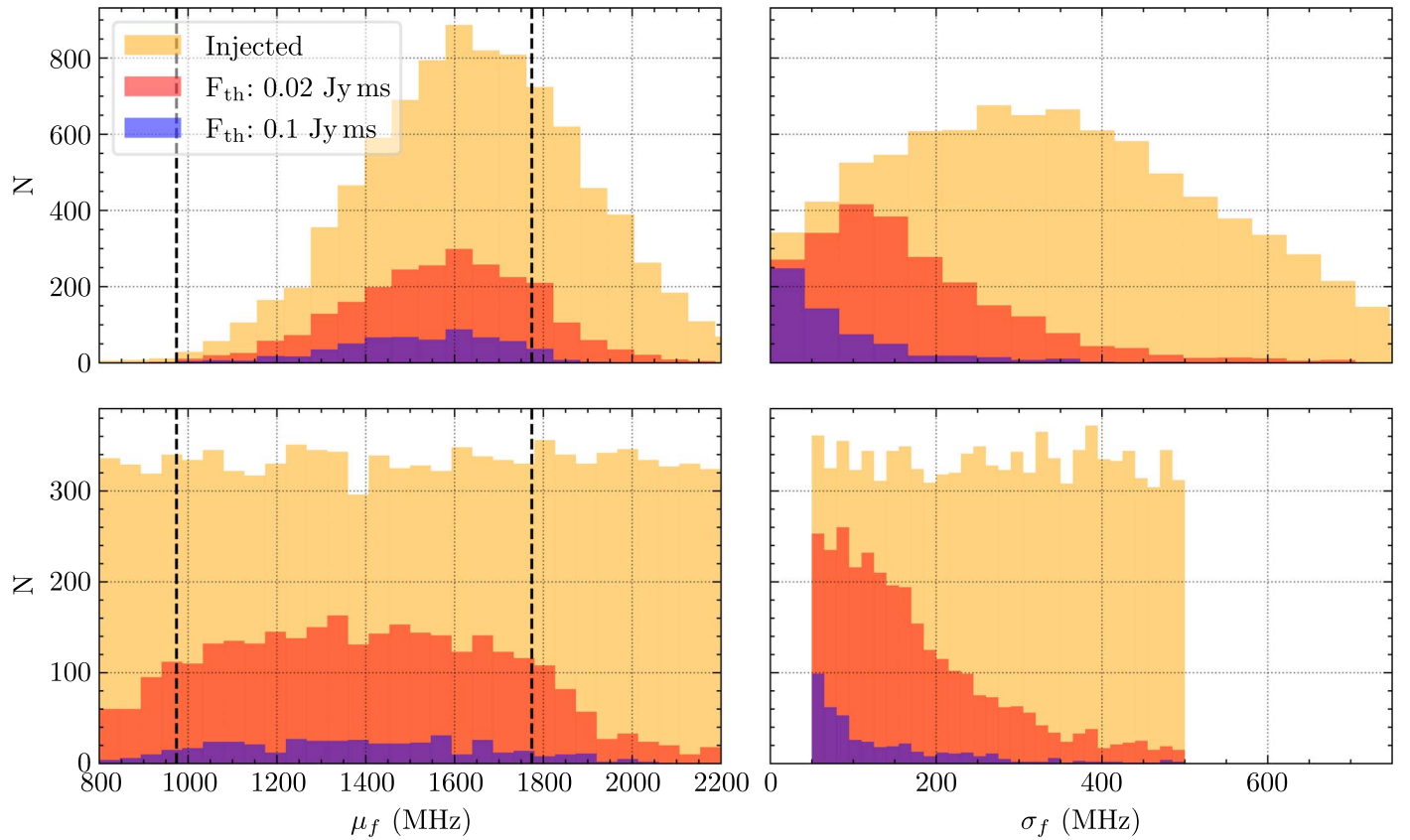
The observed normal distribution of  $\sigma_f$  reported by Aggarwal et al. (2021a) can also be explained only by an intrinsic normal distribution of  $\sigma_f$  and is not recovered by a uniform distribution of  $\sigma_f$ . This can also be seen in the right panels in Figure 3.

### 3.4. Calculating Energy

As mentioned previously, we used Equation (3) to estimate the energy of the burst. This equation uses the burst bandwidth along with the fluence to estimate the energy. This formalism is the way to estimate the burst energy when the emission is not broadband. Under the assumption that emission is broadband, a standard technique to estimate energy from fluence uses center frequency of the band, instead of burst bandwidth (Zhang 2018). This is given by

$$E = 4\pi 10^{-23} \left( \frac{D_L}{\text{cm}} \right)^2 \left( \frac{S}{\text{Jy s}} \right) \left( \frac{\nu_c}{\text{Hz}} \right) \text{erg}, \quad (4)$$

where  $\nu_c$  is the center frequency of the observing band. Recently, Li et al. (2021) used this latter definition of energy and found that the burst energy of FRB 121102 follows a bimodal distribution, using a large sample of bursts. We



**Figure 3.** Distribution of mean ( $\mu_f$ ) and standard deviation ( $\sigma_f$ ) of burst spectra. The intrinsic distribution is shown in yellow, and the distribution of bursts detected at various fluence thresholds is shown in red and blue. The top panels consider a normal intrinsic distribution of  $\mu_f$  and  $\sigma_f$ , while bottom panels assume a uniform intrinsic distribution. Vertical dotted lines mark the observing band. See Section 3.3 for more details.

recalculated the burst energies using the burst bandwidths and fluences reported in their Supplementary Table 1 (Li et al. 2021) and compared them with the energies used by Li et al. (2021).

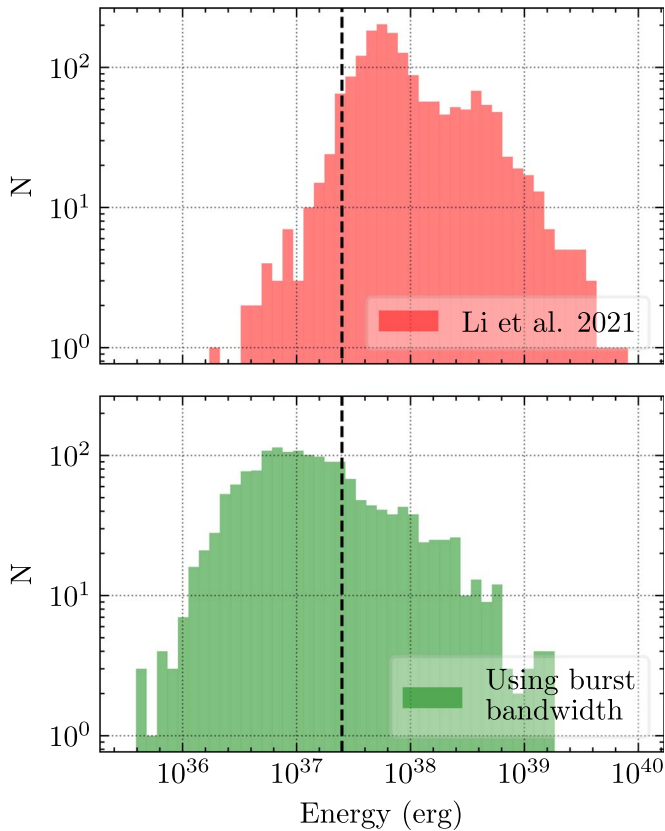
Figure 4 shows the distribution of energies calculated using these two techniques. The distribution of energies does not show any bimodality when burst bandwidths, instead of center frequency, are used to estimate energies. Moreover, the resulting distribution of energies is similar to the ones shown in Figure 2, implying that this result is likely affected by band incompleteness. The method of using center frequency to estimate burst energy makes two key assumptions: (1) emission is broadband, and (2) spectral index is zero, i.e., the emission does not depend on frequency. But the emission from repeaters is characteristically band-limited and Gaussian-like. Therefore, none of these two assumptions are valid for repeating FRBs and so burst bandwidths should be used to estimate energy.

Figure 5 shows the distribution of energies (derived using burst bandwidths) of FRB 121102 bursts detected with FAST and Arecibo (Aggarwal et al. 2021a; Li et al. 2021), along with the recovered distribution of detected bursts using some fiducial values for intrinsic properties similar to the simulation example discussed earlier. Therefore, our analysis shows that the observed shapes in the energy distributions of FRB 121102 can occur due to observational effects caused by bandedness along with a normal distribution of the peak and width of the burst spectra. Notably, both these effects were reported for the bursts presented in Li et al. (2021) and Aggarwal et al. (2021a).

We performed power-law fits on the cumulative distribution of energies (derived using burst bandwidths) from FRB 121102 bursts detected by FAST. We only used the bursts above the energy of  $1.2 \times 10^{37}$  erg, estimated from the 95% completeness limit (0.06 Jy ms) of FAST observations reported by Li et al. (2021) and mean burst bandwidth of FAST bursts (200 MHz). We fitted these bursts using a single power law and a broken power law. The fitted slope for the single power-law fit was  $-0.716 \pm 0.002$ . The slopes (below and above the break energy) for the broken power-law fit were  $-0.693 \pm 0.001$  and  $-1.12 \pm 0.02$  with a fitted break energy of  $(1.05 \pm 0.02) \times 10^{38}$  erg. The high energy slope is consistent with the results of Cruces et al. (2020), while they are inconsistent with those reported by Aggarwal et al. (2021a). It should be noted that the energies of FAST bursts were not corrected for the incompleteness due to the banded nature of bursts. Further, the fluences were derived using S/N and not using fitting. As discussed, both these effects could lead to the incorrect estimation of energy distribution for this sample.

#### 4. Discussion

The results presented in the previous section caution against estimating completeness limits without accounting for the variety of spectral properties of the FRB, especially repeaters. It further complicates the interpretation of energy distributions and properties intrinsic to the FRB. Two main challenges stand out due to the band-limited nature of repeater bursts: (1) detecting the band-limited bursts, and (2) robust estimation of fluence and bandwidth of bursts that are only partially within



**Figure 4.** Energy distribution of FRB 121102 bursts reported by Li et al. (2021). The top panel shows energies calculated using the center frequency of the band (i.e., 1.25 GHz), while the bottom panel shows energies calculated using the bandwidths of the bursts. Vertical dashed line represents the 90% completeness limit of FAST observations estimated by Li et al. (2021). The bottom distribution does not show the bimodality seen in the top figure.

the observing band. In the following subsections, we discuss these two challenges.

#### 4.1. Observing Bandwidth and Subbanded Searches

One possible technique to mitigate the first challenge listed above is to have a large observing bandwidth. A larger observing bandwidth would make it more likely for the burst spectra to fall within the observing band and aid in detecting more bursts.

Another way is to perform subband searches instead of searching full observational bandwidth (R. Anna-Thomas et al. 2021, in preparation; Kumar et al. 2021). The observing band is divided into multiple subbands, on which a single-pulse search is then performed. This strategy is more sensitive to weak band-limited single pulses that will not exceed the detection threshold with a more traditional full-band single-pulse search.

For the example presented in the previous section, we estimated the increase in the number of detected bursts using a subband search as compared to a full-band search. We divided the observing band into four subbands of 200 MHz each and then performed the single-pulse search. We detected more bursts for all three fluence threshold cases (an increase of 5%, 16%, and 41%). More bursts, as expected, reduced the observational biases and resulted in recovered properties that were more representative of the intrinsic distributions.

#### 4.2. Calculating Fluences and Bandwidths

The second challenge listed earlier was a robust estimation of fluence and bandwidth of detected bursts. There is no agreed-upon method to estimate the fluence of the bursts as mentioned previously (see Section 2.2); signal within the observing band is typically used to estimate burst fluence. In some cases, this signal has been modeled using a power law, a Gaussian, or a running power law (Aggarwal et al. 2021a; Pleunis et al. 2021; The CHIME/FRB Collaboration et al. 2021) to estimate the burst fluence and other properties. This is because different FRBs show different spectral properties. In this analysis, we have assumed repeater spectra to be Gaussian, as has been reported by observational campaigns on the most studied repeaters (Law et al. 2017; Pastor-Marazuela et al. 2021; Aggarwal et al. 2021a; Li et al. 2021; Platts et al. 2021). The advantage of assuming a functional form for the spectra is that it enables the estimation of fluence and bandwidth of the burst signal, even if the whole burst is not visible in the band (Aggarwal et al. 2021a). This can provide a robust estimate of the intrinsic fluence and bandwidth of the burst if the assumed functional form is correct.

#### 4.3. Estimating Intrinsic Properties

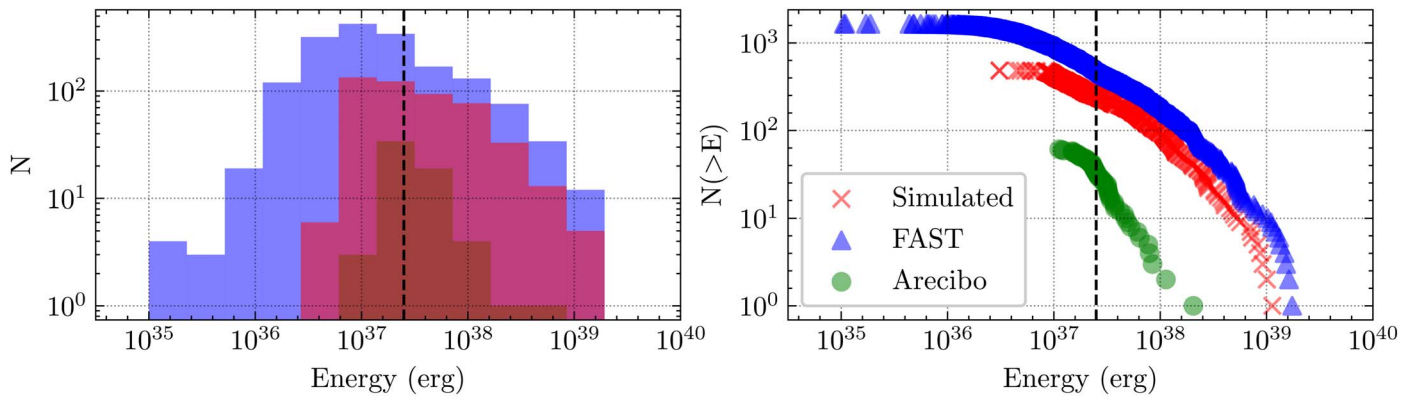
It is now possible to establish a hierarchical framework to infer the intrinsic properties of the repeater bursts. Such a framework would require the following ingredients: (1) observed (or preferably fitted) fluences of the detected bursts from a repeater, (2) completeness limit of the observing system estimated using rigorous injection analysis (Agarwal et al. 2020; Gupta et al. 2021), (3) DM grid (or DM tolerance) used in single-pulse search, and (4) Boxcar widths searched. It would need to assume a distribution for burst energies, a spectral shape, and optionally an intrinsic DM and width distribution.

#### 4.4. Effects of Power-law Slope

We also tested the effect of power-law slope on the energy and spectral parameter distributions of the detected bursts. We reanalyzed the example listed above with two more power-law slopes,  $-1.2$  and  $-1.8$ , and observed the recovered distributions. We did not detect any significant difference in the results for these two power-law slopes with respect to the ones presented earlier, for a slope of  $-1.5$ . The cumulative distribution of energy showed a similar flattening and turnover with decreasing sensitivity. The in-band bursts still provided a more robust and accurate estimate of the intrinsic slope. The burst energy distributions also showed similar shapes as reported in Section 3.2.

### 5. Conclusion

In this Letter, we have discussed various observational effects that can arise due to the banded nature of the spectra of repeating FRBs. Primarily, the banded nature of burst spectra leads to a nonuniform completeness limit across the observing band. This is because many bursts that lie primarily outside the observing band will not be detected. Therefore, contrary to what is typically understood, the search pipeline is not complete to all the fluences above the sensitivity limit determined using traditional injection analysis (Agarwal et al. 2020; Pastor-Marazuela et al. 2021; Gupta et al. 2021). This



**Figure 5.** Energy distributions of repeater bursts. The left panel shows the histogram of energies, and the right panel shows the cumulative distribution of those energies. Blue triangles represent FRB 121102 bursts reported by Li et al. (2021), where we used burst bandwidth to estimate the burst energy. Green circles show the in-band FRB 121102 bursts reported by Aggarwal et al. (2021a). Red crosses show the bursts that were detected using the simulation example discussed in Section 3. Black dashed lines show the 90% completeness limit of FAST observations (Li et al. 2021). The distribution of the observed FRB 121102 energies is very similar to those recovered using simulations. See Section 3.4 for more details.

incompleteness must be accounted for in the analysis that is used to determine the intrinsic properties of the FRB. Also, it is challenging to estimate the intrinsic fluence and bandwidth of bursts that lie on the edge of the observing band. We assumed the spectra to have a Gaussian shape and simulated bursts from an FRB 121102-like source. We then showed that the energy distribution of the detected bursts looks substantially different from the intrinsic distribution and might show peculiar shapes. We also showed that modeling the burst spectra using a Gaussian to determine the intrinsic fluence and bandwidth provides more robust results than traditional approaches.

A normal intrinsic distribution of  $\mu_f$  and  $\sigma_f$  can explain the observed distribution seen for FRB 121102 (Aggarwal et al. 2021a), if the peak of the  $\mu_f$  distribution lies toward the higher frequency end of the observing band. We point out that burst bandwidths, instead of center frequency, should be used to estimate the energy of the banded repeaters from fluence. We also showed that the bimodality in the energy distribution of FRB 121102 bursts reported by Li et al. (2021) disappears when the energy is estimated using burst bandwidths instead of the center frequency of the band. Based on our tests, we can make the following recommendations for single-pulse search and analysis of repeater bursts:

1. Fluence and energies derived using fitting should be preferred over the ones estimated from S/N of the burst. Moreover, only the bursts whose peak and bulk of the emission lie primarily within the observing band (i.e., the in-band bursts) should be used to make inferences about the intrinsic distribution of energies and other properties of the FRB.
2. If the cumulative energy distribution shows a break in the power law, then the higher energy power law could follow the intrinsic distribution. A smooth turnover in the power law will probably not represent the intrinsic distribution in the absence of a break. Even still, if the observations are not very sensitive, it might be impossible to recover the intrinsic properties of the repeater (see Section 3.1). Moreover, these effects depend on the observed fluences from the source. An energetic repeater, which is also close to us, might be easier to interpret than one that is further away.

3. Subbanded searches are more sensitive to such band-limited bursts and will aid in resolving some of the observational biases listed in this Letter.
4. Analysis to determine the search pipeline completeness should incorporate band-limited spectra of FRBs in the simulated FRB injections.

We note that these conclusions apply only to band-limited transient emission, i.e., they might not necessarily apply to situations where the intrinsic burst bandwidth is much greater than the bandwidth of the observing hardware in use. All the analysis scripts and notebooks used in this work are provided in a Github repository.

We would like to thank Devansh Agarwal, Sarah Burke-Spolaor and Casey J. Law for useful discussions and comments on the manuscript. We also thank Di Li for sharing the properties of FRB 121102 bursts detected with FAST. K.A. acknowledges support from NSF grants AAG-1714897 and #2108673.

*Software:* Astropy (Astropy Collaboration et al. 2013; Price-Whelan et al. 2018), Numpy (Harris et al. 2020), Matplotlib (Hunter 2007), Pandas (McKinney 2010; Reback et al. 2021), SciPy (Virtanen et al. 2020), SciencePlots (Garrett 2021).

## ORCID iDs

Kshitij Aggarwal  <https://orcid.org/0000-0002-2059-0525>

## References

- Aggarwal, D., Lorimer, D. R., Surnis, M. P., et al. 2020, *MNRAS*, 497, 352  
 Aggarwal, K., Aggarwal, D., Lewis, E. F., et al. 2021a, arXiv:2107.05658  
 Aggarwal, K., Burke-Spolaor, S., Law, C. J., et al. 2021b, *ApJ*, 914, 53  
 Aggarwal, K., Law, C. J., Burke-Spolaor, S., et al. 2020, *RNAAS*, 4, 94  
 Astropy Collaboration, Robitaille, T. P., Tollerud, E. J., et al. 2013, *A&A*, 558, A33  
 Beniamini, P., & Kumar, P. 2020, *MNRAS*, 498, 651  
 Bera, A., & Chengalur, J. N. 2019, *MNRAS*, 490, L12  
 Cheng, Y., Zhang, G. Q., & Wang, F. Y. 2020, *MNRAS*, 491, 1498  
 Cordes, J. M., Wasserman, I., Hessels, J. W. T., et al. 2017, *ApJ*, 842, 35  
 Cruces, M., Spitler, L. G., Scholz, P., et al. 2020, *MNRAS*, 500, 448  
 Garrett, J. D. 2021, SciencePlots v1.0.9, Zenodo, doi:10.5281/zenodo.4106649  
 Gupta, V., Flynn, C., Farah, W., et al. 2021, *MNRAS*, 501, 2316  
 Harris, C. R., Millman, K. J., van der Walt, S. J., et al. 2020, *Natur*, 585, 357  
 Hessels, J. W. T., Spitler, L. G., Seymour, A. D., et al. 2019, *ApJL*, 876, L23

- Hunter, J. D. 2007, *CSE*, 9, 90
- Keane, E. F., & Petroff, E. 2015, *MNRAS*, 447, 2852
- Kumar, P., Shannon, R. M., Flynn, C., et al. 2021, *MNRAS*, 500, 2525
- Law, C. J., Abruzzo, M. W., Bassa, C. G., et al. 2017, *ApJ*, 850, 76
- Li, D., Wang, P., Zhu, W. W., et al. 2021, arXiv:2107.08205
- Lu, W., & Kumar, P. 2018, *MNRAS*, 477, 2470
- McKinney, W. 2010, in Proc. 9th Python in Science Conf., ed. S. van der Walt & J. Millman (Austin, TX: SciPy), 56
- Metzger, B. D., Margalit, B., & Sironi, L. 2019, *MNRAS*, 485, 4091
- Pastor-Marazuela, I., Connor, L., van Leeuwen, J., et al. 2021, *Natur*, 596, 505
- Petroff, E., Hessels, J. W. T., & Lorimer, D. R. 2019, *A&ARv*, 27, 4
- Platts, E., Caleb, M., Stappers, B. W., et al. 2021, *MNRAS*, 505, 3041
- Pleunis, Z., Good, D. C., Kaspi, V. M., et al. 2021, arXiv:2106.04356
- Price-Whelan, A. M., Sipőcz, B. M., Günther, H. M., et al. 2018, *AJ*, 156, 123
- Reback, J., McKinney, W., Jbrockmendel, et al. 2021, pandas-dev/pandas: Pandas 1.2.1, Zenodo, doi:10.5281/ZENODO.3509134
- Simard, D., & Ravi, V. 2020, *ApJL*, 899, L21
- The CHIME/FRB Collaboration, Amiri, M., Andersen, B. C., et al. 2021, arXiv:2106.04352
- Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, *NatMe*, 17, 261
- Zhang, B. 2018, *ApJL*, 867, L21